



Quantum Computing Using Linear Optics

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Quantum computers are expected to be able to solve mathematical problems that cannot be solved using conventional computers. Many of these problems are of practical importance, especially in the areas of cryptography and secure communications. APL is developing an optical approach to quantum computing in which the bits, or “qubits,” are represented by single photons. Our approach allows the use of ordinary (linear) optical elements that are generally available as off-the-shelf components. Recent experimental demonstrations of a variety of logic gates for single photons, a prototype memory device, and other devices will be described.

INTRODUCTION

In the early 1980s, Richard Feynman demonstrated that there are fundamental limitations in trying to perform simulations of complex quantum systems on conventional computers, regardless of their size or speed.¹ He noted, however, that these problems could be overcome, at least in principle, by building computers based on quantum mechanics instead of classical physics. One naturally wondered if these “quantum computers” would be useful for other applications if they could eventually be built. The answer was shown to be “yes” when, nearly a decade later, Peter Shor discovered a quantum computing algorithm² for efficiently factoring large integers—a problem that has no efficient solution on conventional computers and forms the basis of many secure communications protocols. Soon thereafter, it was shown that a quantum computer could also be used to search an unstructured database much faster than any conventional computer.³ Because of these critical

theoretical developments, there has been a recent explosion of experimental work aimed at building a quantum computer. Researchers in many different areas of physics are actively pursuing a variety of methods to accomplish this challenging goal.

APL is currently developing an optical approach to quantum computing in which the quantum bits, or “qubits,” of information are represented by the quantum state of single photons. For example, the logical value 0 can be represented by a horizontally polarized photon, while the logical value 1 can be represented by a vertically polarized photon. Alternatively, 0 and 1 could be represented by the presence of a single photon in one of two optical fibers (Fig. 1). As described in detail in an earlier *Technical Digest* article,⁴ the fundamental computational advantage of a quantum computer is that quantum mechanics allows the qubits to be in so-called superposition states that do not correspond to specific

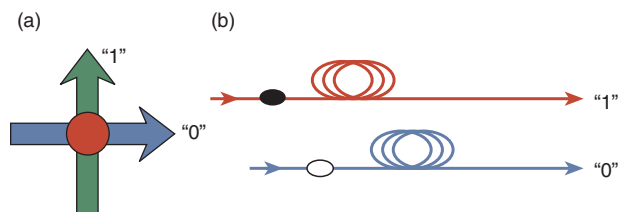


Figure 1. Two methods for implementing quantum bits, or “qubits,” using the quantum states of single photons.⁴ (a) Polarization encoding in which a horizontally polarized single photon represents a logical value of 0 and a vertically polarized single photon represents a logical value of 1. (b) Path encoding, where the presence of a single photon in one of two optical fibers represents a logical value of 0 or 1.

values of 0 or 1. In stark contrast to classical bits (which always have a definite value of either 0 or 1), the qubits can, in some sense, behave as if they had the values of 0 and 1 at the same time.

In addition to photons, many other physical quantum systems are being considered for use as qubits. For example, a single two-level atom in its ground state could correspond to a 0, while the same atom in its excited state would correspond to a 1. Research along these lines is being pursued actively within the context of ion-trap⁵ and nuclear magnetic resonance (NMR)⁶ approaches.

LINEAR OPTICS QUANTUM COMPUTING

The primary advantage of an optical approach to quantum computing is that it would allow quantum logic gates and quantum memory devices to be easily connected using optical fibers or waveguides in analogy with the wires of a conventional computer. This affords a type of modularity that is not readily available in other approaches. For example, the transfer of qubits from one location to another in ion-trap or NMR systems is a very complex process.

The main drawback to an optical approach has been the implementation of the quantum logic gates needed to perform calculations. An important example of a quantum logic gate is the so-called controlled-NOT (CNOT) gate, which has been shown to be a universal gate for quantum computers in the same way that the classical NAND gate is a universal gate for conventional computers.⁷ In other words, any conceivable quantum logic gate can be constructed from a circuit of CNOT gates and single-qubit gates, which are trivial in an optical approach.

As described in Ref. 4, a CNOT gate has two inputs—a control qubit and a target qubit—and operates in such a way that the NOT operation (bit flip) is applied to the target qubit, provided the control qubit has a logical value of 1. Such a logic operation is inherently nonlinear because the state of one quantum particle must be able to control the state of the other. In an optical

approach, this is equivalent to requiring a nonlinear interaction between two single photons, which is typically an extremely weak effect. Conventional nonlinear optical effects, such as frequency doubling of a light beam, are usually only observed in experiments involving intense laser pulses containing billions of photons.⁸ Although several ingenious methods for producing nonlinear interactions at single-photon intensity levels have been considered, they are thought to be either too weak⁹ or accompanied by too much loss¹⁰ to be useful for practical quantum CNOT gates.

It has recently been shown, however, that near-perfect optical quantum logic gates, such as a CNOT gate, can be implemented without the need for a nonlinear interaction between two single photons.¹¹ Logic gates of this kind can be constructed using only linear optical elements such as mirrors and beamsplitters, additional resource photons, and triggering signals from single-photon detectors. In this “linear optics quantum computing” (LOQC) approach, the required nonlinearity arises from the quantum measurement process associated with the detection of the additional resource (“ancilla”) photons.¹¹ Roughly speaking, a single-photon detector either goes off or not, which is a very nonlinear response.

The basic idea of a LOQC-type CNOT gate is illustrated in Fig. 2. Besides the control and target photons,

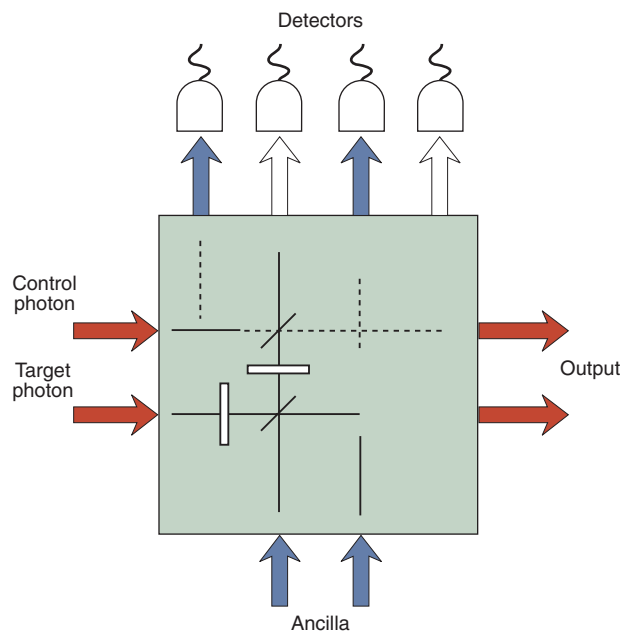


Figure 2. Basic idea of a two-input quantum logic gate constructed using linear optical elements, additional resource (ancilla) photons, and single-photon detectors. The ancilla photons are combined with the logical qubits using linear elements such as beamsplitters and phase shifters. The quantum state of the ancilla photons is measured after they leave the device. The correct logical output is known to have been produced when measurements on the ancilla photons produce certain results. The output can be corrected for other measurement results.

the ancilla photons are injected into a “black box” containing only linear optical elements. The optics are designed so that there are three types of outcome from the device, each signaled by a unique combination of triggering events at a series of single-photon detectors. In one set of outcomes, we know that the control and target photons are in the desired logical output state. In the second type, the control and target photons are known to be in the wrong output state, but they can be corrected in a known way using real-time corrections called “feed-forward control.”¹² The third type of outcome indicates that the control and target photons have been lost or are in a logical state that cannot be corrected.

These LOQC logic gates are referred to as “probabilistic devices” because they occasionally fail, but it is known when a failure has occurred. In addition, the gates can be designed so that the probability of a failure event P_f can be made arbitrarily small. In the original LOQC proposal¹¹ it was shown that P_f can be proportional to $1/N$, where N is the number of ancilla photons consumed by the gate. In a subsequent paper¹³ we described an alternative approach in which P_f scales as $1/N^2$, which greatly reduces the resources required for a given gate fidelity.

EXPERIMENTAL QUANTUM LOGIC GATES

Our goal was to design LOQC logic devices that are as simple, stable, and as robust as possible. We therefore used qubits represented by the polarization states of single photons, as illustrated in Fig. 1a. Polarization-encoded qubits are more resistant to certain kinds of experimental errors and easier to manipulate than the “path-encoded” qubits of Fig. 1b.

The use of polarization-based qubits allowed us to design a CNOT gate using only two polarizing beamsplitters, two polarization-sensitive detectors, and two ancilla photons, as shown in Fig. 3.¹⁴ In this device, the two ancilla photons are in a quantum-mechanically correlated or “entangled” state, where the logical value (i.e., polarization) of each of the ancilla photons is totally undefined but measurements will always find the logical values of the two photons to be the same. Quantum-mechanical correlations of this kind are stronger than those allowed by classical physics, and the possibility of their existence prompted Einstein to question the completeness of quantum mechanics in the early days of the theory.¹⁵ Nonetheless, advances in modern technology have allowed entangled states of this kind to be produced and measured in the laboratory.¹⁶

The device shown in Fig. 3 exploits the entanglement of the ancilla pair to implement the desired CNOT logic operation on the input control and target qubits. The operation of this gate requires that the control photon

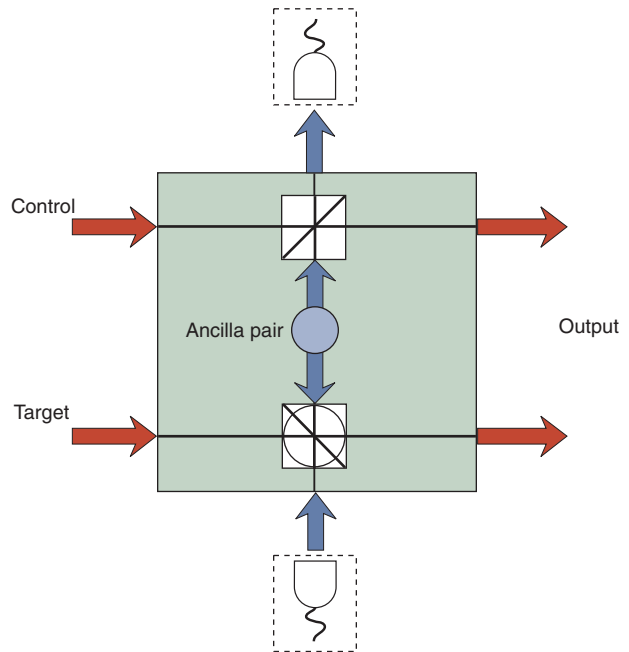


Figure 3. Overview of the APL linear optics quantum controlled-NOT gate.¹⁴ In addition to the input control and target qubits, the device consists of two polarizing beamsplitters, an entangled pair of resource ancilla photons, and two single-photon detectors. The output is known to be correct when each detector registers one photon, which occurs with a probability of 25%.

and one member of the entangled ancilla photon simultaneously arrive at the upper-polarizing beamsplitter, while the target photon and second member of the entangled ancilla pair simultaneously arrive at the lower-polarizing beamsplitter. The correct logical output is known to have been produced whenever each of the detectors registers one and only one photon, which occurs with a probability of 25%.¹⁴

We recently demonstrated a CNOT gate of the kind shown in Fig. 3.¹⁷ In our experiment, the arrangement of the polarizing beamsplitters was altered so that the role of the entangled photon pair could be replaced by a single ancilla photon propagating through the entire device. This simplified the technical requirements of the experiment by reducing the total number of photons involved from four to three. However, in this simplified configuration the operation of the logic gate could only be verified by measuring (and thus destroying) the control and target qubits after they exited the device. Nonetheless, the experimental results represented the first demonstration of a CNOT gate for single photons and a tangible step toward full-scale quantum computing using linear optics. Figure 4 is a photograph of the experimental apparatus.

In this proof-of-principle experiment, one of the three required photons was obtained from an attenuated laser pulse, while the other two were produced through the process of spontaneous parametric down-conversion (SPDC).¹⁸ The SPDC process involves passing an

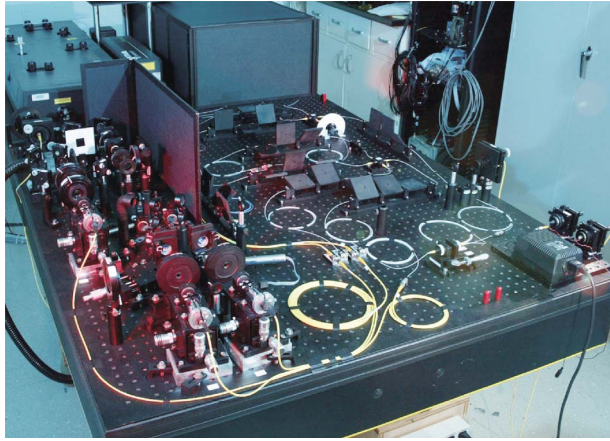


Figure 4. Photograph of the experimental apparatus used to demonstrate the first quantum controlled-NOT logic gate for single-photon qubits. The experimental components included a mode-locked Ti-Sa laser, single-mode fiber components, parametric down-conversion photon sources, and low-noise single-photon detectors.

intense laser pumping pulse through a nonlinear medium so that correlated pairs of photons occasionally emerge. SPDC is a purely quantum-mechanical phenomenon that can be viewed as the annihilation of a single photon in the laser pulse followed by the creation of two lower-energy down-conversion photons, under the conditions that energy and momentum are conserved.

The operation of our CNOT logic gate relied on multiphoton quantum interference effects that required the three photons to be indistinguishable, aside from their polarizations (i.e., logical values). This required a combination of precise spectral filtering, the use of single-mode optical fibers for spatial mode-matching, and timing precision on the order 10^{-13} s. Once these parameters were optimized, the CNOT gate could be tested using polarizing optics to control the values of the control and target input qubits and by using polarization analyzers followed by single-photon detectors to test and measure the output of the device.

An example of these types of measurements is shown in Fig. 5. In comparison with the theory, the experimental results clearly demonstrate the desired logical truth table of a CNOT gate, aside from technical errors on the order of 10%. Experiments aimed at reducing these technical errors and using an entangled ancilla pair of resource photons as shown in Fig. 3 are currently under way at APL.

In addition to a CNOT gate, we have also demonstrated a number of other important LOQC gates including a quantum parity check, an exclusive-OR gate,¹⁹ and

a quantum encoding device²⁰ capable of encoding (copying) the value of a single qubit into a logical state represented by two photons. This encoding operation can be used to provide redundancy that can protect against photon loss or other errors which may occur in a realistic environment. In addition to its use in quantum computing applications, an encoder of this kind can be employed to implement a “quantum relay” that would help extend the range of quantum cryptography systems.²¹

SINGLE-PHOTON SOURCES

All of the logic gates described above rely on small numbers of ancilla photons and typically operate with probabilities of success in the range of 25 to 50%. Although these gates have a number of important applications, full-scale quantum computers will require highly efficient gates involving larger numbers of ancilla photons.²² Therefore, the development of reliable sources capable of emitting a single photon at well-defined time intervals is a key ingredient in the realization of LOQC.

The development of a source of single photons is a critical requirement and one that cannot be accomplished by re-engineering a conventional light source. For example, the number of photons in a laser pulse follows a Poisson distribution. Attenuating a train of laser pulses until each pulse contains an average of one photon will result in a few pulses that actually contain zero photons or more than one. Injecting those pulses into an LOQC logic gate would produce undesirable errors. Although quantum error correction techniques that are analogous to classical parity checks do exist, it is necessary to keep the intrinsic error rate below a certain threshold on the order of 1%.

One method for realizing a true single-photon source is through spontaneous emission from an isolated

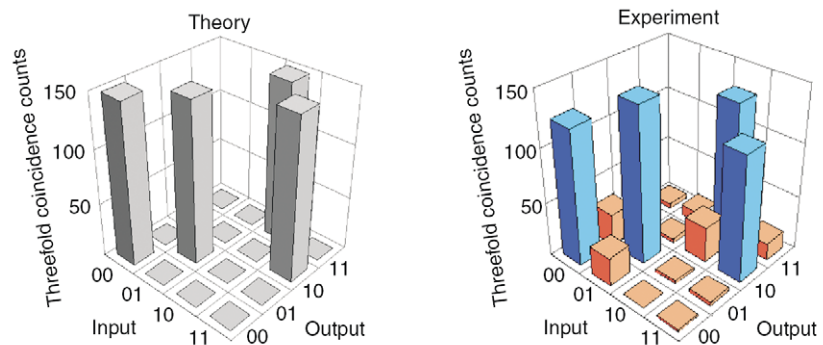


Figure 5. Experimental results demonstrating the logical truth table for the first quantum controlled-NOT gate for single photons. The NOT operation (bit flip) is applied to the target qubit if and only if the control qubit has the logical value of 1. The experimental data are seen to be in agreement with theory, aside from average technical errors on the order of 10%. (Reprinted, with permission, from Ref. 17, Fig. 3; © 2003 by the American Physical Society.)

two-level quantum system. For example, a single two-level atom in its excited state can emit only one photon, as was experimentally verified nearly 30 years ago.²³ In recent years, the suitability of a variety of other single-photon emitters has been investigated, including single molecules, quantum dots, and solid-state defects such as color centers.²⁴

In contrast to these approaches, APL's contribution has been the development of a single-photon source based on the photon pairs produced in SPDC. The basic idea of this source is illustrated in Fig. 6. Because SPDC is known to produce pairs of photons, the detection of one photon of a pair can be used to signal the presence of the twin photon. A high-speed optical switch is then used to store the twin photon in a storage loop until it is needed, at which time it can be switched back out of the storage loop. We have experimentally demonstrated a single-photon source of this kind. Its performance is currently limited by losses in the optical switch, but we are now developing a low-loss switch for single photons. Storage loops and low-loss switches can also form the basis of a quantum memory device for single-photon qubits.²⁷

QUANTUM CIRCUITS

As we mentioned earlier, one of the main advantages of an optical approach to quantum computing is the ability to connect logic and memory devices using optical fibers in analogy with the use of wires in conventional electronic circuits. To demonstrate this capability, we recently constructed and tested a relatively simple quantum circuit that combines two linear optics quantum logic gates²⁸ to perform a useful function. The circuit consisted of two probabilistic exclusive-OR (XOR) gates in series, as shown in Fig. 7. Two single-photon qubits formed the input to the first XOR gate. The output of that gate then served as the input to the second XOR gate, along with a third single-photon qubit.

Because an XOR gate can be used to measure the parity of its inputs, it can be shown that the circuit of Fig. 7 calculates the parity of the three input qubits when

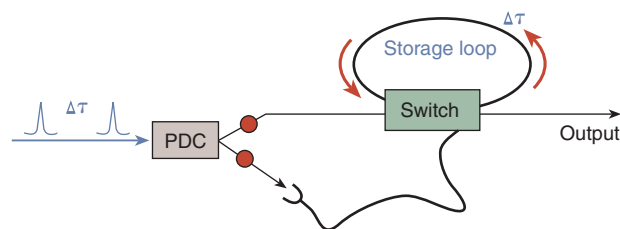


Figure 6. Overview of APL's single-photon source.²⁵ A parametric down-conversion crystal (PDC) is pumped by a train of laser pulses separated in time by $\Delta\tau$, which causes it to randomly emit pairs of correlated photons.¹⁸ Once a pair is emitted, the detection of one of the photons activates an electro-optic switch that is used to route the other photon into a storage loop. The stored photon can then be switched back out of the loop when it is needed. (Reprinted, with permission, from Ref. 26.)

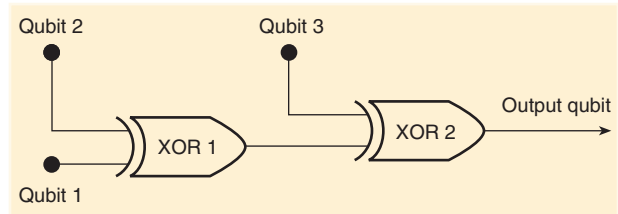


Figure 7. A simple logic circuit for photonic qubits. This circuit calculates the parity of three input qubits using two exclusive-OR (XOR) logic gates. We have demonstrated the classical truth table for this circuit along with a variety of quantum interference effects when the input qubits are in quantum superposition states.

they all have specific logical values of 0 or 1. When one or more of the input qubits is in a more general superposition state, the circuit of Fig. 7 produces a quantum mechanical output state that cannot be reproduced by any classical device. We have demonstrated both the classical truth table corresponding to well-defined input values as well as a variety of quantum interference effects associated with superposition states. Although XOR gates are not reversible, this simple circuit is still useful in a number of important applications and demonstrates our ability to connect independent devices using optical fibers to form more complex circuits.

SUMMARY

In the past few years, the prospects for quantum computing have progressed from a fascinating academic exercise to one with important applications. A linear optics approach appears to be a promising method for eventually building a full-scale quantum computer. APL has been actively developing and demonstrating many of the basic building blocks that are required, including the first CNOT gate for single-photon qubits¹⁷ and a prototype single-photon source,²⁵ a quantum memory device,²⁷ and a photon number resolving detector.²⁹ In addition, we recently demonstrated the first quantum circuit for photonic qubits.²⁸

Although these proof-of-principle experiments are encouraging, many significant technical challenges remain. In particular, the need for large numbers of ancilla to achieve low error rates may increase the complexity of such a quantum computer and increase the amount of resources required. With that in mind, we are investigating a variety of ways to reduce the dependence on large numbers of ancilla photons. For example, we have recently shown that the quantum Zeno effect can be used to completely suppress the failure events that would otherwise occur in these probabilistic quantum logic gates.³⁰ This new approach would eliminate the need for ancilla photons and make LOQC more practical for large-scale applications. We are planning to perform experiments of that kind in the near future.

In summary, the development of quantum computers would have a major impact in a number of areas of

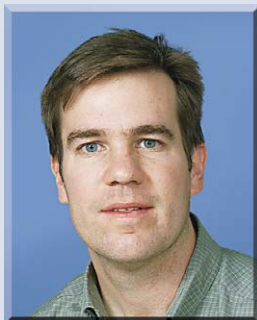
practical importance, including cryptography, secure communications, and optimization problems. APL has demonstrated the basic building blocks of an optical approach to quantum computing. Although a number of challenges remain, our approach appears promising as a method for eventually building a full-scale quantum computer.

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REFERENCES

- ¹Feynman, R. P., "Quantum Mechanical Computers," *Opt. News* **11**, 11–20 (1985).
- ²Shor, P. W., "Algorithms for Quantum Computation: Discrete Logarithms and Factoring" in *35th Ann. Symp. on Foundations of Computer Science: Proc.*, S. Goldwasser (ed.), IEEE Computer Society Press (1994).
- ³Grover, L. K., "Quantum Mechanics Helps in Searching for a Needle in a Haystack," *Phys. Rev. Lett.* **79**, 325–329 (1997).
- ⁴Franson, J. D., and Jacobs, B. C., "Quantum Computing," *Johns Hopkins APL Tech. Dig.* **18**(2), 188–192 (1997).
- ⁵Monroe, C., Meekhof, D. M., King, B. E., Itano, W., and Wineland, D. J., "Demonstration of a Fundamental Quantum Logic Gate," *Phys. Rev. Lett.* **75**, 4714–4717 (1995).
- ⁶Cory, D. G., Fahmy, A. F., and Havel, T. F., "Ensemble Quantum Computing by NMR Spectroscopy," *Proc. Nat. Acad. Sci. USA* **94**, 1634–1639 (1997).
- ⁷Nielsen, M. A., and Chuang, I. L., *Quantum Computation and Quantum Information*, Cambridge University Press, Cambridge, UK (2000).
- ⁸Shen, Y. R., *The Principles of Nonlinear Optics*, Wiley Interscience, New York (1984).
- ⁹Turchette, Q. A., Hood, C. J., Lange, W., Mabuchi, H., and Kimble, H. J., "Measurement of Conditional Phase Shifts for Quantum Logic," *Phys. Rev. Lett.* **75**, 4710–4713 (1995).
- ¹⁰Franson, J. D., "Cooperative Enhancement of Optical Quantum Gates," *Phys. Rev. Lett.* **78**, 3852–3855 (1997).
- ¹¹Knill, E., LaFlamme, R., and Milburn, G. J., "A Scheme for Efficient Quantum Computation with Linear Optics," *Nature* **409**, 46–52 (2001).
- ¹²Pittman, T. B., Jacobs, B. C., and Franson, J. D., "Demonstration of Feed-Forward Control for Linear Optics Quantum Computation," *Phys. Rev. A* **66**, 052305 (2002).
- ¹³Franson, J. D., Donegan, M. M., Fitch, M. J., Jacobs, B. C., and Pittman, T. B., "High Fidelity Quantum Logic Operations Using Linear Optical Elements," *Phys. Rev. Lett.* **89**, 137901 (2004).
- ¹⁴Pittman, T. B., Jacobs, B. C., and Franson, J. D., "Probabilistic Quantum Logic Operations Using Polarizing Beam Splitters," *Phys. Rev. A* **64**, 062311 (2001).
- ¹⁵Einstein, A., Podolsky, B., and Rosen, N., "Can Quantum Mechanical Description of Physical Reality Be Considered Complete?" *Phys. Rev.* **48**, 696–702 (1935).
- ¹⁶Pittman, T. B., and Franson, J. D., "Violation of Bell's Inequality with Photons from Independent Sources," *Phys. Rev. Lett.* **90**, 24041 (2003).
- ¹⁷Pittman, T. B., Fitch, M. J., Jacobs, B. C., and Franson, J. D., "Experimental Controlled-NOT Logic Gate for Single Photons in the Coincidence Basis," *Phys. Rev. A* **68**, 032316-3 (2003).
- ¹⁸Klyshko, D. N., *Photons and Nonlinear Optics*, Gordon and Breach Science, New York (1988).
- ¹⁹Pittman, T. B., Jacobs, B. C., and Franson, J. D., "Demonstration of Non-Deterministic Quantum Logic Operations Using Linear Optical Elements," *Phys. Rev. Lett.* **88**, 257902 (2002).
- ²⁰Pittman, T. B., Jacobs, B. C., and Franson, J. D., "Probabilistic Quantum Encoder for Single-Photon Qubits," *Phys. Rev. A* **69**, 042306 (2004).
- ²¹Jacobs, B. C., Pittman, T. B., and Franson, J. D., "Quantum Relays and Noise Suppression Using Linear Optics," *Phys. Rev. A* **66**, 052307 (2002).
- ²²Franson, J. D., Donegan, M. M., and Jacobs, B. C., "Generation of Entangled Ancilla States for Use in Linear Optics Quantum Computing," *Phys. Rev. A* **66**, 052328 (2004).
- ²³Kimble, H. J., Dagenis, M., and Mandel, L., "Photon Antibunching in Resonance Fluorescence," *Phys. Rev. Lett.* **39**, 691–694 (1977).
- ²⁴Kurtsiefer, C., Mayer, S., Zarda, P., and Weinfurter, H., "Solid State Source of Single Photons," *Phys. Rev. Lett.* **85**, 290–293 (2000).
- ²⁵Pittman, T. B., Jacobs, B. C., and Franson, J. D., "Single Photons on Pseudodemand from Stored Parametric Down-Conversion," *Phys. Rev. A* **66**, 042303 (2002).
- ²⁶Pittman, T. B., Fitch, M. J., Jacobs, B. C., and Franson, J. D., "Periodic Single-Photon Source of Quantum Memory," in *Proc. SPIE* **5161**, pp. 57–65 (2003).
- ²⁷Pittman, T. B., and Franson, J. D., "Cyclical Quantum Memory for Photonic Qubits," *Phys. Rev. A* **66**, 062302 (2002).
- ²⁸Pittman, T. B., Jacobs, B. C., and Franson, J. D., "Experimental Demonstration of Quantum Circuit Using Linear Optics Quantum Gates," *Phys. Rev. A* (2004). (e-print archive quant-ph/0404059).
- ²⁹Fitch, M. J., Jacobs, B. C., Pittman, T. B., and Franson, J. D., "Photon-number Resolution Using Time-Multiplexed Single-Photon Detectors," *Phys. Rev. A* **68**, 043814 (2003).
- ³⁰Franson, J. D., Jacobs, B. C., and Pittman, T. B., "Quantum Computing Using Single Photons and the Zeno Effect," *Phys. Rev. A*. (2004). (e-print archive quant-ph/0401133).

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